

The goal of this chapter is to prove the following theorem:

Theorem (Anseny, 1975):

Assume that $d \geq 2$ and let $f^{in} \in L^1 \cap L^\infty(\mathbb{R}^d_x \times \mathbb{R}^d_v)$ satisfy $f^{in} \geq 0$ and $\mathcal{E}^{in} < \infty$. Then there exists a global weak solution $f \in L^\infty(\mathbb{R}_+; L^1(\mathbb{R}^d_x \times \mathbb{R}^d_v))$ of the Vlasov-Poisson system with initial data f^{in} . Moreover, this solution satisfies

$$0 \leq f(t, x, v) \leq \|f^{in}\|_{L^\infty}$$

for a.e. $(x, v) \in \mathbb{R}^d_x \times \mathbb{R}^d_v$ and for all $t \geq 0$, together with mass bound

$$\int_{\mathbb{R}^d_x \times \mathbb{R}^d_v} f(t, x, v) dx dv \leq \mathcal{M}^{in} < \infty$$

for all $t \geq 0$ and the energy bound

$$\int_{\mathbb{R}^d_x \times \mathbb{R}^d_v} \frac{1}{2} |v|^2 f(t, x, v) dx dv + \int_{\mathbb{R}^d_x} \frac{1}{2} |\nabla_x \Phi(t, x)|^2 dx \leq \mathcal{E}^{in} < \infty$$

for a.e. $t \geq 0$.

The initial condition is verified in the sense of distributions, i.e. for all $\varphi \in C_c^\infty(\mathbb{R}^d \times \mathbb{R}^d)$, the function

$t \mapsto \int_{\mathbb{R}^d \times \mathbb{R}^d} f(t, x, v) \varphi(x, v) dx dv$
is continuous on \mathbb{R}_+ and satisfies

$$\int_{\mathbb{R}^d \times \mathbb{R}^d} f(t, x, v) \varphi(x, v) dx dv = \int_{\mathbb{R}^d \times \mathbb{R}^d} f_{in}(x, v) \varphi(x, v) dx dv.$$

Recall that the Cauchy problem for the Vlasov-Poisson system is

$$\begin{cases} \partial_t f(t, x, v) + v \cdot \nabla_x f(t, x, v) - \nabla_x \Phi(t, x) \cdot \nabla_v f(t, x, v) = 0 \\ \text{in } \mathbb{R}_+^1 \times \mathbb{R}^d \times \mathbb{R}^d, \\ -\Delta_x \Phi(t, x) = \int_{\mathbb{R}^d} f(t, x, v) dv, \\ f(0, x, v) = f_{in}(x, v). \end{cases}$$

To prove this theorem from Arsénauy, we first introduce the approximate Vlasov-Poisson system (this is done via mollifiers).

We introduce mollifiers as follows.
 Let $\zeta \in C^\infty(\mathbb{R}^d)$ satisfy $\zeta(x) = \zeta(-x)$
 ≥ 0 for all $x \in \mathbb{R}^d$, $\text{supp}(\zeta) \subset \mathbb{B}(0, 1)$,
 and $\int_{\mathbb{R}^d} \zeta(x) dx = 1$, and set

$$\zeta_\epsilon(x) = \epsilon^{-d} \zeta(x/\epsilon),$$

$$\zeta_\epsilon(u, v) = \zeta_\epsilon(x) \zeta(v).$$

We use even mollifiers so that
 the convolution operator is
 self-adjoint on $L^2(\mathbb{R}^N)$.

Lemma:

Let $\chi \in C^\infty(\mathbb{R}^N)$ satisfy $\chi(x) = \chi(-x)$
 for all $x \in \mathbb{R}^N$. Then the convolution
 operator $C_\chi: \phi \mapsto \chi * \phi$ is
 self-adjoint on $L^2(\mathbb{R}^N)$.

Proof:

First, we show that C_χ is bounded.
 Let $\phi \in L^2(\mathbb{R}^N)$. Then

$$\|C_\chi \phi\|_{L^2} \leq \|\chi\|_{L^1} \|\phi\|_{L^2}$$

by Young's inequality for the
 convolution. Hence to show that
 it is self-adjoint, it is sufficient
 to verify that it is symmetric.
 Let $\psi \in L^2(\mathbb{R}^N)$. Then

$$\begin{aligned} (C_\chi \phi, \psi) &= (\chi * \phi, \psi) \\ &= (\phi, \chi * \psi). \end{aligned}$$

Indeed, by Fubini's Theorem and since χ is even,

$$\begin{aligned}
 \int_{\mathbb{R}^N} \chi * \phi(x) \psi(x) dx &= \int_{\mathbb{R}^N} \left(\int_{\mathbb{R}^N} \chi(x-t) \phi(t) dt \right) \psi(x) dx \\
 &= \int_{\mathbb{R}^N} \phi(t) \left(\int_{\mathbb{R}^N} \chi(t-x) \psi(x) dx \right) dt \\
 &= \int_{\mathbb{R}^N} \phi(t) \chi * \psi(t) dt.
 \end{aligned}$$

We define the approx. Helmholtz-Poisson system for a given $\epsilon > 0$ as \square

$$\left\{ \begin{array}{l}
 \partial_t \phi_\epsilon + v \cdot \nabla_x \phi_\epsilon - \nabla_x \Phi_\epsilon \cdot \nabla_v \phi_\epsilon = 0, \\
 -\Delta_x \Phi_\epsilon(t, \cdot) = \xi_\epsilon * \xi_\epsilon * \rho_\epsilon(t, \cdot), \\
 \nabla_x \Phi_\epsilon \rightarrow 0 \text{ as } |x| \rightarrow \infty, \\
 \rho_\epsilon(t, x) = \int_{\mathbb{R}^d} f(t, x, v) dv, \\
 \phi_\epsilon|_{t=0} = \xi_\epsilon * (\mathbf{1}_{\{|x| < 1\}} \mathbf{1}_{\{|v| < 1\}})^{\text{in}} =: \phi^{\text{in}}.
 \end{array} \right. \quad (\text{PV}_\epsilon)$$

Thus, for all $t \geq 0$,

$$-\nabla_x \Phi_\epsilon(t, \cdot) = (\xi_\epsilon * \xi_\epsilon * \nabla G_d) * \rho_\epsilon(t, \cdot)$$

where G_d is the d -dimensional Green's function, the fundamental solution.

For each $\epsilon > 0$, one has

$$\zeta_\epsilon * \zeta_\epsilon * \nabla \zeta_\epsilon \in C^\infty(\mathbb{R}^d) \cap L^{d, \infty}(\mathbb{R}^d) \cap L^\infty(\mathbb{R}^d)$$

Proposition:

For each $f^{\text{in}} \in L^1 \cap L^\infty(\mathbb{R}^d \times \mathbb{R}^d)$ satisfying $f^{\text{in}} \geq 0$ and

$$\int_{\mathbb{R}^d \times \mathbb{R}^d} \frac{1}{2} |v|^2 f_\epsilon^{\text{in}}(x, v) dx + \int_{\mathbb{R}^d} \frac{1}{2} |\tilde{E}_\epsilon^{\text{in}}(x)|^2 dx = E^{\text{in}} < \infty,$$

with

$$\tilde{E}_\epsilon^{\text{in}} = -\zeta_\epsilon * \nabla \zeta_\epsilon * \rho_\epsilon^{\text{in}}$$

where

$$\rho_\epsilon^{\text{in}}(x) = \int_{\mathbb{R}^d} f_\epsilon^{\text{in}}(x, v) dv,$$

there exists a unique weak solution $f_\epsilon \in C(\mathbb{R}_+; L^1(\mathbb{R}^d \times \mathbb{R}^d))$ of (VP_ϵ) .
This solution satisfies

$$0 \leq f_\epsilon(t, x, v) \leq \|f^{\text{in}}\|_{L^\infty(\mathbb{R}^d \times \mathbb{R}^d)}$$

for all $(x, v) \in \mathbb{R}^d \times \mathbb{R}^d$ and $t \geq 0$,
together with mass conservation

$$\int_{\mathbb{R}^d \times \mathbb{R}^d} f_\epsilon(t, x, v) dx dv = \int_{\mathbb{R}^d \times \mathbb{R}^d} f^{\text{in}}(x, v) dx dv$$

and the approximate energy conservation bound

$$\int_{\mathbb{R}^d \times \mathbb{R}^d} \frac{1}{2} |v|^2 f_\varepsilon(t, x, v) dx dv + \int_{\mathbb{R}^d} \frac{1}{2} |\tilde{E}_\varepsilon(t, x)|^2 dx \leq \varepsilon^{in}$$

for all $t \geq 0$, where

$$\tilde{E}_\varepsilon(t, \cdot) = (\zeta_\varepsilon * \nabla \zeta_\varepsilon) * \rho_\varepsilon(t, \cdot).$$

Proof:

Wlog, we may assume that $\int_{\mathbb{R}^d \times \mathbb{R}^d} f_\varepsilon^{in}(x, v) dx dv = 1$ and that $\int_{\mathbb{R}^d \times \mathbb{R}^d} v f_\varepsilon^{in}(x, v) dx dv = 0$. Since the dynamics of the approximate Vlasov-Poisson system preserves the total mass and total momentum, one has

$$\int_{\mathbb{R}^d \times \mathbb{R}^d} f_\varepsilon(t, x, v) dx dv = 1 \text{ for all } t \geq 0,$$

and

$$\int_{\mathbb{R}^d \times \mathbb{R}^d} v f_\varepsilon(t, x, v) dx dv = 0 \text{ for all } t \geq 0.$$

We have already seen that the Vlasov-Poisson system can be put in the form

$$\partial_t f(t, z) + \operatorname{div}_z (f(t, z) \int_{\mathbb{R}^d} K(z, z') f(t, z') dz') = 0,$$

with $z = (x, v)$ and

$$K_\varepsilon(x, v, x', v') = (v - v', \xi_\varepsilon * \xi_\varepsilon * \nabla \psi).$$

We have seen that in this formalism, the existence and uniqueness is obtained via the local solution of the mean field flow. Indeed, there exists a unique map

$$\begin{aligned} \mathbb{R}_+ \times \mathbb{R}^{2d} \times \mathcal{P}_1(\mathbb{R}^d) &\rightarrow \mathbb{R}^{2d} \\ (t, z^{\text{in}}, \mu^{\text{in}}) &\mapsto Z_\varepsilon(t, z^{\text{in}}, \mu^{\text{in}}) \end{aligned}$$

such that $t \mapsto Z_\varepsilon(t, z^{\text{in}}, \mu^{\text{in}})$ is the integral curve of the vector field

$$\begin{aligned} z &\mapsto \int_{\mathbb{R}^{2d}} K_\varepsilon(z, z') \mu_\varepsilon(z') dz' \\ &:= (K \mu_\varepsilon(t))(z) \end{aligned}$$

passing through z^{in} at time $t=0$, where $\mu_\varepsilon(t) := Z_\varepsilon(t, \cdot, \mu^{\text{in}}) \# \mu^{\text{in}}$.
Set

$$Z_\varepsilon(t, z^{\text{in}}) := Z_\varepsilon(t, z^{\text{in}}, (\varepsilon \mathcal{L}^{2d}))$$

and

$$V_\varepsilon(t, z) := (K \mu_\varepsilon(t))(z)$$

where $\mu_\varepsilon(t) := Z_\varepsilon(t, \cdot, (\varepsilon \mathcal{L}^{2d}))$.

We have that

$$V_\varepsilon(t, z) = (v - v', -(\xi_\varepsilon * \xi_\varepsilon * \nabla G_d) * \rho_\varepsilon).$$

Hence, $V_\varepsilon \in C(\mathbb{R}_+ \times \mathbb{R}^{2d}; \mathbb{R}^{2d})$ and $V_\varepsilon(\cdot, \cdot) \in C^\infty(\mathbb{R}^{2d}; \mathbb{R}^{2d})$. Moreover, we have an estimate on $V_\varepsilon(t, z)$. By Young's inequality, we get

$$\begin{aligned} |V_\varepsilon(t, z)| &\leq |v - v'| + \|\rho_\varepsilon\|_{L^1} \|\xi_\varepsilon * \xi_\varepsilon * \nabla G_d\|_{L^\infty} \\ &\leq |z| + \|\xi_\varepsilon\|_{L^1} \|\xi_\varepsilon * \nabla G_d\|_{L^\infty} \\ &\leq |z| + \|\xi_\varepsilon * \nabla G_d\|_{L^\infty}. \end{aligned}$$

Hence $Z_\varepsilon \in C^1(\mathbb{R}_+ \times \mathbb{R}^{2d}; \mathbb{R}^{2d})$ and $Z_\varepsilon(\cdot, \cdot)$ is a C^∞ -diffeomorphism of \mathbb{R}^{2d} for all $t \geq 0$. Since $V_\varepsilon(t, \cdot)$ is divergence free, the diffeomorphism $Z_\varepsilon(\cdot, \cdot)$ leaves the Lebesgue measure \mathbb{Z}^{2d} invariant. Therefore, the solution of (VPE) is of the form

$$\begin{aligned} \mu(t) &= Z_\varepsilon(\cdot, \cdot) \# (\rho_\varepsilon^{\text{in}} \mathbb{Z}^{2d}) \\ &= \rho_\varepsilon^{\text{in}} (Z_\varepsilon(\cdot, \cdot)^{-1}) \mathbb{Z}^{2d} \\ &:= \rho_\varepsilon(t, \cdot) \mathbb{Z}^{2d} \end{aligned}$$

for all $t \geq 0$.

Then mass conservation is satisfied since $\tilde{z}_\epsilon(t, \cdot)$ is a diffeomorphism leaving invariant the Lebesgue measure. Continuity is obvious. We have to check the energy conservation. Since $f_\epsilon \in C_c^\infty(\mathbb{R}^d \times \mathbb{R}^d)$, we have that $f_\epsilon \in C^1(\mathbb{R}_+ \times \mathbb{R}^{2d})$ and $\text{supp}(f_\epsilon(t, \cdot))$ is compact for all $t \geq 0$. By Lebesgue's dominated convergence theorem we get

$$\begin{aligned}
 & \frac{d}{dt} \int_{\mathbb{R}^{2d}} \frac{1}{2} |V|^2 f_\epsilon(t, x, v) \, d(x, v) \\
 &= \int_{\mathbb{R}^{2d}} \frac{1}{2} |V|^2 \partial_t f_\epsilon(t, x, v) \, d(x, v) \\
 &= \int_{\mathbb{R}^{2d}} \frac{1}{2} |V|^2 (\nabla_x \Phi_\epsilon(t, x) \cdot \nabla_v f_\epsilon(t, x, v) \\
 &\quad - v \cdot \nabla_x f_\epsilon(t, x, v)) \, d(x, v) \\
 &= \int_{\mathbb{R}^{2d}} -\text{div}_x (\frac{1}{2} |V|^2 v f_\epsilon(t, x, v)) \, d(x, v) \\
 &\quad + \int_{\mathbb{R}^{2d}} \frac{1}{2} |V|^2 \nabla_v f_\epsilon(t, x, v) \cdot \nabla_x \Phi_\epsilon(t, x) \, d(x, v).
 \end{aligned}$$

The first term vanishes using the divergence theorem and the fact that f_ϵ has compact support. By integration by parts, we get

$$\begin{aligned}
 \frac{d}{dt} \int_{\mathbb{R}^{2d}} |v|^2 f_\epsilon(t, x, v) dx dv & \\
 &= - \int_{\mathbb{R}^{2d}} v \cdot \nabla_x \Phi_\epsilon(t, x) f_\epsilon(t, x, v) dx dv \\
 &= - \int_{\mathbb{R}^d} j_\epsilon(t, x) \cdot \nabla_x \Phi_\epsilon(t, x) dx \\
 &= \int_{\mathbb{R}^d} j_\epsilon(t, x) \cdot (\xi_\epsilon * \tilde{E}_\epsilon(t, \cdot)) dx \\
 &= \int_{\mathbb{R}^{2d}} (\xi_\epsilon * j_\epsilon(t, \cdot))(x) \cdot \tilde{E}_\epsilon(t, x) dx
 \end{aligned}$$

for all $t \geq 0$ where we use the fact that the convolution operator with even functions is self-adjoint for the last equality.

Then, using Green's Identity,
we get

$$\frac{d}{dt} \int_{\mathbb{R}^{2d}} \frac{1}{2} |v|^2 f_\epsilon(y, v) dy dv$$

$$= \int_{\mathbb{R}^d} (y_\epsilon * \operatorname{div}_x \partial_t f_\epsilon(\cdot, \cdot))(x) \tilde{\Phi}_\epsilon(y, x) dx$$

$$= - \int_{\mathbb{R}^d} (y_\epsilon * \partial_t (p_\epsilon(\cdot, \cdot)))(x) \tilde{\Phi}_\epsilon(y, x) dx$$

$$= \int_{\mathbb{R}^d} (\partial_t \Delta_x \tilde{\Phi}_\epsilon(y, x)) \tilde{\Phi}_\epsilon(y, x) dx$$

$$= - \int_{\mathbb{R}^d} \partial_t \nabla_x \tilde{\Phi}_\epsilon(y, x) \cdot \nabla_x \tilde{\Phi}_\epsilon(y, x) dx$$

$$= - \frac{d}{dt} \int_{\mathbb{R}^d} \frac{1}{2} |\nabla_x \tilde{\Phi}_\epsilon(y, x)|^2 dx$$

for all $t \geq 0$, and hence the
energy is conserved. □