

Separation of variables - reloaded

Exercise 1: how to handle Robin BCs.

Solve, by sep. of variables, the IBVP

$$\begin{cases}
 0 < x < 1 \\
 t > 0
 \end{cases}
 \left\{
 \begin{array}{l}
 u_t = \alpha^2 u_{xx} \\
 u(0, t) = 0 \\
 u_x(1, t) + \ell u(1, t) = 0 \\
 u(x, 0) = x
 \end{array}
 \right.$$

physical groundwater > 0

Solution: the solution is given by an (infinite) sum of special solutions of the form $X(x)T(t)$.

So let's plug in this $\xrightarrow{\text{Ansatz and solve}}$ for X and T .

Step 1: $T'(t)X(x) = \alpha^2 X''(x)T(t)$

$$\rightarrow \underbrace{\frac{X''(x)}{X(x)}}_{\mu} = \underbrace{\frac{T'(t)}{\alpha^2 T(t)}}_{\mu} \quad (\mu \in \mathbb{R})$$

\rightarrow both constant

We then look at the equation for $X(x)$ together w/ BCs.

$$\left\{
 \begin{array}{l}
 X(0)T(t) = 0 \\
 X'(1)T(t) + \ell X(1)T(t) = 0
 \end{array}
 \right.$$

\rightsquigarrow non-trivial solutions

$$\left\{
 \begin{array}{l}
 X(0) = 0 \\
 X'(1) + \ell X(1) = 0
 \end{array}
 \right.$$

So we get for $X(x)$ the problem (*)

$$\left\{
 \begin{array}{l}
 X'' = \mu X \\
 X(0) = 0 \\
 X'(1) + \ell X(1) = 0
 \end{array}
 \right.$$

(This is a special case of a Sturm-Liouville pb.)

Step 2: discussing the separation constant.

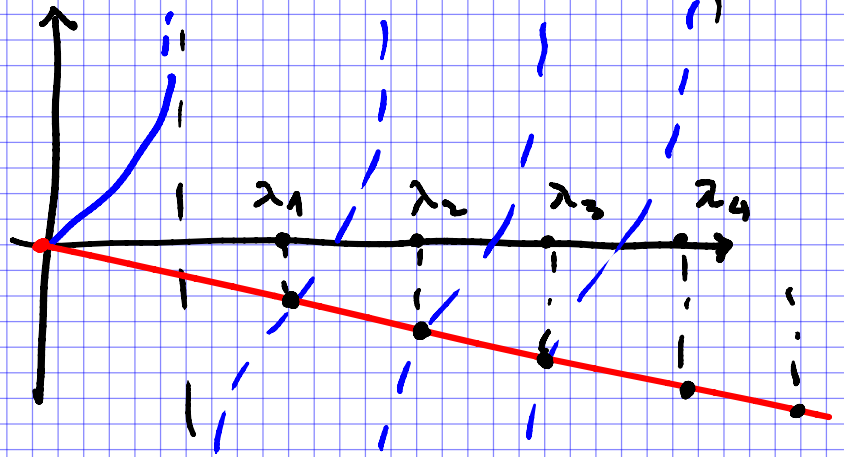
distinguish 3 cases based on sign(μ):

• $\mu = 0 \implies x(x) = A + Bx \quad (\forall A, B \in \mathbb{R})$
 $x(0) = 0 = A \implies A = 0$
 $x'(1) + kx(1) = B + kB = 0$
 $= B(1+k) \implies B = 0$

• $\mu > 0 \implies x(x) = A \cosh(\lambda x) + B \sinh(\lambda x) \quad (\forall A, B \in \mathbb{R})$
 $(\mu = \lambda^2)$
 $\lambda > 0$
 $x(0) = 0 = A \implies A = 0$
 $x'(1) + kx(1) = 0 = \lambda B \cosh(\lambda) + k B \sinh(\lambda)$
 $= B \lambda \cosh(\lambda) \left(1 + k \frac{\tanh(\lambda)}{\lambda} \right)$
positive if $\lambda > 0$
 $\implies B = 0$

• $\mu < 0 \implies x(x) = A \cos(\lambda x) + B \sin(\lambda x) \quad (\forall A, B \in \mathbb{R})$
 $\mu = -\lambda^2$
 $(\lambda > 0)$
 $x(0) = 0 = A \implies A = 0$
 $x'(1) + kx(1) = B \lambda \cos(\lambda) + k B \sin(\lambda)$
 $= B \lambda \cos(\lambda) \left(1 + k \frac{\tan(\lambda)}{\lambda} \right)$
 $= 0$

so if we want non-trivial solutions we must impose
 $\tan \lambda = -\frac{1}{k} \lambda$



we get a sequence $\{\lambda_n\} \quad n \geq 1$

so $X_n(x) = \sin(\lambda_n x) \quad n \geq 1$, and correspondingly

$T_n(t) = e^{-(\alpha \lambda_n)^2 t}$ and so our Ansatz is that

$$u(x, t) = \sum_{n \geq 1} a_n e^{-(\alpha \lambda_n)^2 t} \sin(\lambda_n x)$$

and we must determine a_1, a_2, \dots so to match the I.C, i.e.

$$\sum_{n \geq 1} a_n \sin(\lambda_n x) = x$$

Step 3: to find a_m apply "brute force trick"

multiply by $\sin(\lambda_m x)$ and integrate

$$\int_0^1 x \sin(\lambda_m x) dx = \sum_{n \geq 1} a_n \int_0^1 \sin(\lambda_n x) \sin(\lambda_m x) dx$$

$$= \begin{cases} 0 & \text{if } n \neq m \\ \frac{\lambda_m - \sin \lambda_m \cos \lambda_m}{2 \lambda_m} & \text{if } n = m \end{cases}$$

$$\int_0^1 x \sin(\lambda_m x) dx = a_m \left(\frac{\lambda_m - \sin \lambda_m \cos \lambda_m}{2 \lambda_m} \right)$$

$$a_m = \frac{2 \lambda_m}{\lambda_m - \sin \lambda_m \cos \lambda_m} \int_0^1 x \sin(\lambda_m x) dx$$

To be discussing convergence of series $\sum_{n \geq 1} a_n \sin(\lambda_n x)$ (see below) the problem is solved.

Probk. (1) in general, if we start from

$$\begin{cases} u_t = \alpha^2 u_{xx} \\ \alpha_1 u_x(0, t) + \beta_1 u(0, t) = 0 \\ \alpha_2 u_x(1, t) + \beta_2 u(1, t) = 0 \\ u(x, 0) = \phi(x) \end{cases}$$

one can argue as above, w/ (*) is substituted by the Sturm-Liouville ph.

$$\begin{cases} X'' = \mu X \\ \alpha_1 X'(0) + \beta_1 X(0) = 0 \quad k_1(t) \\ \alpha_2 X'(1) + \beta_2 X(1) = 0 \quad k_2(t) \end{cases}$$

eigenvalues of $\frac{d^2}{dx^2}$ under oblique BCs

by general "abstract" results there always exist

sequences of eigenvalues $\lambda_1 \leq \lambda_2 \leq \lambda_3 \dots$
 eigenfunctions $x_1 \ x_2 \ x_3 \dots$

w/ $\int_0^1 x_i(x) x_j(x) dx = 0$ if $i \neq j$

∴ (pointwise) convergence theorems hold true

(2) Standard trick to reduce to homogeneous BCs
 "by subtracting an offset function". Discussed in detail in Fowlow's Lesson 6, and problem 5.1.

Exercise 2 : solve inhomogeneous PDEs

$0 < x < 1$
 $t > 0$

$$\begin{cases} u_t - \alpha^2 u_{xx} = f(x,t) \\ \alpha_1 u_x(0,t) + \beta_1 u(0,t) = 0 \\ \alpha_2 u_x(1,t) + \beta_2 u(1,t) = 0 \\ u(x,0) = \phi(x) \end{cases}$$

Step 1 (usual story ...) get to SL problem

Step 2 ($\mu = -\lambda^2$)

$$\begin{cases} X'' = \mu X \\ \alpha_1 X'(0) + \beta_1 X(0) = 0 \\ \alpha_2 X'(1) + \beta_2 X(1) = 0 \end{cases}$$

(theory) $\lambda_1 \leq \lambda_2 \leq \dots$
 $x_1 \leq x_2 \leq \dots$ ← these are not explicit in general

w/ $\int_0^1 x_i(x) x_j(x) dx = 0$ if $i \neq j$

Step 3 . we'll write $\phi(x) = \sum_{n \geq 1} a_n x_n(x)$

∴ similarly we decompose $f(x,t) = \sum_{n \geq 1} f_n(t) x_n(x)$
 (eigenfunction expansion)

As always, you compute a_u and $f_u(t)$ by brute force trick. Recall:

$$f(x, t) = \sum_{u \geq 1} f_u(t) X_u(x)$$

$$\int_0^1 f(x, t) X_w(x) dx = \sum_{u \geq 1} f_u(t) \int_0^1 X_u(x) X_w(x) dx$$

orthogonality \rightarrow

$$= f_w(t) \cdot \int_0^1 X_w^2(x) dx$$

$$f_w(t) = \frac{\int_0^1 f(x, t) X_w(x) dx}{\int_0^1 X_w^2(x) dx}$$

Step 4: compute the response to each mode of the source.

We make the usual Ansatz that $u(x, t) = \sum_{u \geq 1} X_u(x) T_u(t)$

where $u_u(x, t) := X_u(x) T_u(t)$ solves the problem

$$\begin{cases} (u_u)_t - \alpha^2 (u_u)_{xx} = f_u(t) X_u(x) \\ \text{usual BCs} \leftarrow \text{already ok} \\ u_u(x, 0) = a_u X_u \end{cases}$$

$$\begin{cases} T_u' X_u - \alpha^2 T_u X_u'' = f_u(t) X_u(x) \\ T_u(0) X_u = a_u X_u \end{cases} \quad \overline{u_u(x, t) = X_u(x) T_u(t)}$$

$$\begin{cases} T_u' + \alpha^2 \lambda_u^2 T_u = f_u(t) \\ T_u(0) = a_u \end{cases}$$

$$\overline{X_u'' = -\lambda_u^2 X_u}$$

how to solve? $\left. \begin{array}{l} y'(t) + a(t)y(t) = b(t) \\ y(t_0) = y_0 \end{array} \right\}$ (Cauchy pl.)
for 1st order ODE

unique sol. $y(t) = e^{-A(t)} \left[y_0 + \int_{t_0}^t e^{A(s)} b(s) ds \right]$
 $A(t) = \int_{t_0}^t a(s) ds$

So in the special case t_0 is question we get

$$T_u(t) = e^{-\alpha^2 \lambda_u^2 t} \left[a_u + \int_0^t e^{\alpha^2 \lambda_u^2 s} f_u(s) ds \right]$$

end of the exercise.

Now, we specify the discussion above to a "concrete case".

$$\begin{cases} u_t - \alpha^2 u_{xx} = \sin(3\pi x) \\ u(0, t) = 0 \\ u(1, t) = 0 \\ u(x, 0) = \sin(\pi x) \end{cases}$$

Step 1 (...)

$$\begin{cases} x'' = \mu x \\ x(0) = 0 \\ x(1) = 0 \end{cases}$$

cf. lecture 3
 Dirichlet BC
 lead to $\sin(-)$
 so standard
 Fourier series

already seen

$\mu = -\lambda^2$
 (Step 2)

$$\lambda_n = (n\pi), \quad n \geq 1$$

$$X_n(x) = \sin(n\pi x)$$

Step 3: decompose source and initial datum w.r.t. basis given in step 1

$$f(x, t) = \sin(3\pi x) = \sum_{n \geq 1} \underbrace{f_n(t)} X_n(x)$$

$$\rightsquigarrow f_3(t) \equiv 1$$

$$f_n(t) \equiv 0 \quad \text{if } n \neq 3$$

$$\phi(x) = \sin(\pi x) = \sum_{n \geq 1} a_n X_n(x)$$

$$\rightsquigarrow \begin{cases} a_1 = 1 \\ a_n = 0 \end{cases}$$

if $n \neq 1$

Step 4: response to each mode

$$(\#) \begin{cases} T_u' + (\alpha\pi u)^2 T_u = f_u(t) \\ T_u(0) = a_u \end{cases}$$

① if $u \neq 1, 3$ then $(\#)$ becomes

$$\begin{cases} T_u'(t) + (\alpha\pi u)^2 T_u = 0 \\ T_u(0) = 0 \end{cases}$$

$$\longrightarrow T_u \equiv 0$$

② if $u = 1$ then $(\#)$ becomes

$$\begin{cases} T_1'(t) + (\alpha\pi)^2 T_1 = 0 \\ T_1(0) = 1 \end{cases}$$

$$\longrightarrow T_1(t) = e^{-(\alpha\pi)^2 t}$$

③ if $u = 3$ then $(\#)$ becomes

$$\begin{cases} T_3'(t) + (3\alpha\pi)^2 T_3 = 1 \\ T_3(0) = 0 \end{cases}$$

$$\longrightarrow T_3(t) = e^{-(3\alpha\pi)^2 t} \int_0^t e^{(3\alpha\pi)^2 s} ds$$

$$\longrightarrow T_3(t) = \frac{1}{(3\alpha\pi)^2} \left[1 - e^{-(3\alpha\pi)^2 t} \right]$$

Now, we put all pieces together. The solution:

$$\begin{aligned} u(x, t) &= X_1(x) T_1(t) + X_3(x) T_3(t) \\ &= e^{-(\alpha\pi)^2 t} \sin(\pi x) + \frac{1}{(3\alpha\pi)^2} \left[1 - e^{-(3\alpha\pi)^2 t} \right] \sin(3\pi x) \end{aligned}$$

Asymptotic behaviour for $t \rightarrow \infty$

$$u(x, t) \xrightarrow{\sim} \frac{1}{(3\alpha\pi)^2} \sin(3\pi x)$$