

$$(1) \quad \check{v}_i^* (v_{s_0} \wedge v_{s_1} \wedge \dots) = \check{v}_i^* (v_{s_0}) v_{s_1} \wedge \dots \\ - \check{v}_i^* (v_{s_1}) v_{s_0} \wedge \dots \\ \uparrow + \dots$$

all summands = 0 unless $s_k = i$ for some k

in which case

$$= \pm v_{s_0} \wedge v_{s_1} \wedge \dots \wedge v_{s_{k+1}} \wedge v_{s_{k-1}} \wedge \dots$$

where the sign is determined by

$$k = (-1)^k v_{s_0} \wedge v_{s_1} \wedge \dots \wedge v_{s_{k+1}} \wedge v_{s_{k-1}} \wedge \dots$$

Similarly for $\hat{v}_i (v_{s_0} \wedge v_{s_1} \wedge \dots)$

$$= v_i \wedge v_{s_0} \wedge v_{s_1} \wedge \dots$$

By moving v_i to its rightfull position
we obtain that

$$\hat{v}_i (v_{s_0} \wedge v_{s_1} \wedge \dots) = (-1)^{k-1} v_{s_0} \wedge \dots \wedge v_{s_k} \wedge v_i \wedge v_{s_{k-1}} \wedge \dots$$

where $s_{k-1} < i < s_k$

and = 0 if $i = s_k$.

(2)

If $s' \neq s \cup \{i\}$, then

$$\langle \hat{v}_i \cdot \psi_s | \psi_{s'} \rangle = 0 = \langle \psi_s | \check{v}_i^* \psi_{s'} \rangle$$

where ψ_s and $\psi_{s'}$ are semi-infinite

monomials associated to sets S & S'

Assume $\hat{S} = S \cup \{i\}$, then

$$\hat{v}_i \cdot \psi_S = (-1)^{k-1} \psi_{S'}$$

such that $s_{k-1} < i < s_k$

$$\text{while } \tilde{v}_i^* \psi_{S'} = (-1)^{k-1} \psi_S$$

$$\Rightarrow \langle \hat{v}_i \psi_S | \psi_{S'} \rangle = \langle \psi_S | \tilde{v}_i^* \psi_{S'} \rangle = (-1)^{k-1}$$

(3) (3.1) Since $\psi_S = v_{s_0} \wedge v_{s_1} \wedge v_{s_2} \wedge \dots$
 $\tilde{v}_m^* \cdot \psi_S = 0$ for all $m > s_0$.

$$\begin{aligned} \Rightarrow \psi_{n_i}^* \cdot \psi_0 &= \left(\sum_{m \geq 0} b_{i,m} \tilde{v}_{m+i}^* \right) \cdot \psi_0 \\ &= \left(\sum_{\substack{m \geq 0 \\ m+i \leq 0}} b_{i,m} \tilde{v}_{m+i}^* \right) \cdot \psi_0 \end{aligned}$$

Note that the sum above involves only finitely many summands

Similarly $\hat{v}_m \cdot \psi_S = 0$ for all $m \leq 0$ because $s_i = i \forall i \leq 0$

$$\begin{aligned} \Rightarrow \psi_{n_i} \cdot \psi_S &= \left(\sum_{m \geq 0} a_{i,m} \hat{v}_{n_i-m} \right) \cdot \psi_S \\ &= \left(\sum_{\substack{m \geq 0 \\ n_i-m \geq k}} a_{i,m} \hat{v}_{n_i-m} \right) \cdot \psi_S \end{aligned}$$

$\uparrow$ $\uparrow$ since without
 again finite sum

$\Rightarrow \langle \psi_0 | \psi_{n_i} \psi_{m_i}^* \psi_0 \rangle$ is well-defined.

$$(3.2) \quad \langle \psi_0 | \psi_{n_i} \psi_{m_i}^* \psi_0 \rangle$$

$$= \langle \psi_{n_i}^* \psi_0 | \psi_{m_i}^* \psi_0 \rangle$$

$\uparrow \hat{v}_i$ and $\hat{v}_i^*$ are adjoints.

By the arguments from (3.1),

$$\hat{v}_i^* n_i - n_i \psi_0 = 0 \quad \text{for } n_i - n > 0$$

$$\& \quad \hat{v}_i^* m_i + m \psi_0 = 0 \quad \text{for } m_i + m > 0$$

$\Rightarrow$ we may assume that

$$a_{i,n} = 0 \quad \text{if } n_i - n > 0$$

*

$$b_{i,m} = 0 \quad \text{if } m_i + m > 0$$

Moreover, by the commutation

relation from the Exercise sheet 6

we have

$$** \quad \psi_{n_i} \psi_{m_j}^* = -\psi_{m_j}^* \psi_{n_i} + \sum_{\substack{n,m \\ m_j+n \\ = n_i-n}} a_{i,n} \cdot b_{j,m} \cdot \text{id}$$

$$\text{and } \heartsuit = \langle \psi_0 | \psi_{n_i} \psi_{m_j}^* \psi_0 \rangle$$

because

$$= \langle \psi_{n_i}^* \psi_0 | \psi_{m_j}^* \psi_0 \rangle$$

and vectors in $\psi_{n_i}^* \psi_0$ pair

non-trivially with vectors in $\psi_{n_j}^* \psi_0$
 iff $m_j + n = n_j - n$.

Let's now prove Wick's theorem
 by induction.

The base case $k=1$ is trivially true.
 Assume Wick's theorem holds for $k-1$.

$$\begin{aligned} \langle \psi_0 | \psi_{n_1} \dots \psi_{n_k} \psi_{m_k}^* \dots \psi_{m_1}^* \psi_0 \rangle &= \\ \langle \psi_0 | \psi_{n_k} \psi_{m_k}^* \psi_0 \rangle \langle \psi_0 | \psi_{n_1} \dots \psi_{n_{k-1}} \psi_{m_{k-1}}^* \dots \psi_{m_1}^* \psi_0 \rangle & \\ - \langle \psi_0 | \psi_{n_1} \dots \psi_{n_{k-1}} \psi_{m_k}^* \psi_{n_k} \psi_{m_{k-1}}^* \dots \psi_{m_1}^* \psi_0 \rangle & \end{aligned}$$

where we used $**$

continue to slide $\psi_{m_k}^*$ to the
 left to obtain

$$\begin{aligned} \sum_{j=0}^{k-1} (-1)^j \langle \psi_0 | \psi_{n_{k-j}} \psi_{m_k}^* \psi_0 \rangle & \\ \cdot \langle \psi_0 | \psi_{n_1} \dots \widehat{\psi_{n_j}} \dots \psi_{n_k} \psi_{m_{k-1}}^* \dots \psi_{m_1}^* \psi_0 \rangle & \\ + (-1)^k \langle \psi_0 | \psi_{n_k}^* \psi_{n_1} \dots \psi_{n_k} \psi_{m_{k-1}}^* \dots \psi_{m_1}^* \psi_0 \rangle & \\ \uparrow & \\ = (-1)^k \langle \psi_{n_k} \psi_0 | \dots \rangle = 0 & \\ \text{by the adjointness} & \quad \text{by } *. \end{aligned}$$

By the induction hypothesis,

$$\begin{aligned}
& \langle \psi_0 | \psi_{n_1} \dots \psi_{n_{k-1}} \dots \psi_{n_k} \psi_{n_{k-1}}^* \dots \psi_{n_1}^* | \psi_0 \rangle \\
&= \det_{\substack{i, 1=1, \dots, k \\ i \neq k, 1 \neq k-1}} \langle \psi_0 | \psi_{n_i} \psi_{n_i}^* | \psi_0 \rangle.
\end{aligned}$$

Combining this with ~~***~~ we

obtain the claim by noticing that

~~***~~ is the expansion of the determinant with respect to a line.

(4)

(4.1) Since M has only finitely many non-zero entries

$\exp(M)$ is defined as in the finite dimensional case.

(4.2) We may assume we are in the finite-dimensional case, in which case it is true (at least for complex matrices) by taking $\log$ of a matrix.

(5) Again, we may assume that we are in the finite-dimensional case, then the claim follows from $\uparrow$ (because μ ...)

expanding

(sets non-trivially
on finitely many v_i)

$$(1 + M + \frac{M^2}{2!} + \dots) v_{s_0} \wedge (1 + M + \frac{M^2}{2!} + \dots) v_{s_{-1}} \wedge \dots \wedge (1 + M + \frac{M^2}{2!} + \dots) v_{s_{-k}}$$

and collecting terms of the
same degree;

$$\begin{aligned} & v_{s_0} \wedge v_{s_{-1}} \wedge \dots \\ + & M v_{s_0} \wedge v_{s_{-1}} \wedge \dots + v_{s_0} \wedge M v_{s_{-1}} \wedge \dots \\ + & M v_{s_0} \wedge M v_{s_{-1}} \wedge v_{s_{-2}} \wedge \dots + v_{s_0} \wedge M v_{s_{-1}} \wedge M v_{s_{-2}} \wedge \dots \\ + & \frac{M^2}{2} v_{s_0} \wedge v_{s_{-1}} \wedge v_{s_{-2}} + v_{s_0} \wedge \frac{M^2}{2} v_{s_{-1}} \wedge v_{s_{-2}} \\ + & \text{terms of order 3} \\ + & \dots \dots \dots \\ = & (1 + M + \frac{M^2}{2} + \dots) \cdot v_{s_0} \wedge v_{s_{-1}} \wedge \dots \end{aligned}$$